

$$I = \int d^4x \sqrt{-g} R$$

Make variation w.r.t.  $g_{\mu\nu}$ , keeping the  $\frac{g_{\mu\nu}}{\partial}$  and their 1<sup>st</sup> derivatives constant on the boundary.

Claim:

Putting  $\delta I = 0$  for arbitrary  $\delta g_{\mu\nu}$  gives Einstein's vacuum equations.

$$(i) R_{\mu\nu} = \partial_\nu \overset{\alpha}{\Gamma}_{\mu\sigma} - \partial_\sigma \overset{\alpha}{\Gamma}_{\mu\nu} - \overset{\alpha}{\Gamma}_{\mu\nu} \overset{\beta}{\Gamma}_{\alpha\beta} + \overset{\alpha}{\Gamma}_{\mu\beta} \overset{\beta}{\Gamma}_{\alpha\nu}$$

$$\Rightarrow R = g^{\mu\nu} R_{\mu\nu}$$

$$= \underbrace{g^{\mu\nu} (\partial_\nu \overset{\sigma}{\Gamma}_{\mu\sigma} - \partial_\sigma \overset{\sigma}{\Gamma}_{\mu\nu})}_{R^*} - \underbrace{g^{\mu\nu} (\overset{\sigma}{\Gamma}_{\mu\nu} \overset{\rho}{\Gamma}_{\sigma\rho} - \overset{\rho}{\Gamma}_{\mu\sigma} \overset{\sigma}{\Gamma}_{\rho\nu})}_L$$

involving 2<sup>nd</sup> derivatives

But, they occur only linearly, so they can be removed by partial integration.

$$\Rightarrow R^* \sqrt{-g} = \overbrace{\partial_\nu (g^{\mu\nu} \overset{\sigma}{\Gamma}_{\mu\sigma} \sqrt{-g}) - \partial_\sigma (g^{\mu\nu} \overset{\sigma}{\Gamma}_{\mu\nu} \sqrt{-g})}^{\text{Total differential} \dots \text{contribute nothing to } I} - \partial_\nu (g^{\mu\nu} \sqrt{-g}) \overset{\sigma}{\Gamma}_{\mu\sigma} + \partial_\sigma (g^{\mu\nu} \sqrt{-g}) \overset{\sigma}{\Gamma}_{\mu\nu}$$

$$\left[ \begin{array}{l} \partial_\nu (g^{\mu\nu} \sqrt{-g}) = -g^{\nu\beta} \overset{\mu}{\Gamma}_{\beta\nu} \sqrt{-g} \\ \text{and } \partial_\sigma (g^{\mu\nu} \sqrt{-g}) = (-g^{\nu\beta} \overset{\mu}{\Gamma}_{\beta\sigma} - g^{\mu\alpha} \overset{\nu}{\Gamma}_{\alpha\sigma} + g^{\mu\nu} \overset{\beta}{\Gamma}_{\sigma\beta}) \sqrt{-g} \end{array} \right]$$

$$\Rightarrow \int d^4x R^* \sqrt{-g} = \int d^4x (g^{\nu\beta} \overset{\mu}{\Gamma}_{\beta\nu} \overset{\sigma}{\Gamma}_{\mu\sigma} \sqrt{-g} - 2g^{\nu\beta} \overset{\mu}{\Gamma}_{\beta\sigma} \overset{\sigma}{\Gamma}_{\mu\nu} \sqrt{-g} + g^{\mu\nu} \overset{\beta}{\Gamma}_{\sigma\beta} \overset{\sigma}{\Gamma}_{\mu\nu} \sqrt{-g})$$

$$\Rightarrow \int d^4x R^* \sqrt{g} = \int -2\sqrt{-g} d^4x$$

$$\Rightarrow I = \int L \sqrt{-g} d^4x$$

which involves only the  $g_{\mu\nu}$  and their first derivatives.

Put  $L = \mathcal{L} \sqrt{-g}$ .

We take it as the action density for the gravitational field.

It's not a scalar density.

But it's more convenient than  $R\sqrt{-g}$ , which is a scalar density, because it does not involve 2<sup>nd</sup> derivatives of the  $g_{\mu\nu}$ .

According to the ordinary ideas of dynamics, the action is the time interval integral of the Lagrangian.

We have

$$I = \int \mathcal{L} d^4x$$

$$= \int x_0 \int \mathcal{L} dx^1 dx^2 dx^3$$

So the Lagrangian is evidently

$$\int \mathcal{L} dx^1 dx^2 dx^3$$

Thus  $L$  may be considered as Lagrangian density (in 3D) as well as the action density (in 4D). We may look upon the  $g_{\mu\nu}$  as dynamical coordinates and their time derivatives as the velocities.

Then the Lagrangian is quadratic (nonhomogeneous) in the velocities, as it usually is in ordinary dynamics.

Vary  $L$  with using  $\Gamma_{\nu\mu}^{\mu} \sqrt{-g} = \partial_{\nu} \sqrt{-g}$

$$\therefore \text{(ii)} \quad \partial_{\nu} g = \frac{\partial g}{\partial \sqrt{-g}} \partial_{\nu} \sqrt{-g} \quad \begin{matrix} \partial_{\nu} \sqrt{-g} = \frac{1}{2} \frac{1}{\sqrt{-g}} \partial_{\nu} g \\ = \frac{1}{2} \frac{\partial g}{\sqrt{-g}} \partial_{\nu} g \end{matrix}$$

$$\text{and } \partial_{\nu} g = g^{\lambda\mu} \partial_{\nu} g_{\lambda\mu}$$

$$\Rightarrow g^{\lambda\mu} \partial_{\nu} (g_{\lambda\mu}) \sqrt{-g} = 2 \sqrt{-g} \partial_{\nu} \sqrt{-g}$$

$$\Rightarrow \boxed{\partial_{\nu} \sqrt{-g} = \frac{1}{2} \sqrt{-g} g^{\lambda\mu} \partial_{\nu} g_{\lambda\mu}}$$

$$\text{(iii)} \quad \Gamma_{\nu\mu}^{\mu} = g^{\lambda\mu} \Gamma_{\lambda\nu\mu} = \frac{1}{2} g^{\lambda\mu} (\partial_{\mu} g_{\lambda\nu} + \partial_{\nu} g_{\lambda\mu} - \partial_{\lambda} g_{\mu\nu})$$

$$= \frac{1}{2} g^{\lambda\mu} \partial_{\nu} (g_{\lambda\mu})$$

$$\Rightarrow \Gamma_{\nu\mu}^{\mu} \sqrt{-g} = \frac{1}{2} \sqrt{-g} g^{\lambda\mu} \partial_{\nu} g_{\lambda\mu} = \partial_{\nu} \sqrt{-g}$$

$$\delta(\Gamma_{\mu\nu}^{\alpha} \Gamma_{\alpha\beta}^{\beta} g^{\mu\nu} \sqrt{g})$$

$$= \Gamma_{\mu\nu}^{\alpha} \delta(\Gamma_{\alpha\beta}^{\beta} g^{\mu\nu} \sqrt{g}) + \Gamma_{\alpha\beta}^{\beta} g^{\mu\nu} \sqrt{g} \delta \Gamma_{\mu\nu}^{\alpha}$$

$$= \Gamma_{\mu\nu}^{\alpha} \delta(g^{\mu\nu} \partial_{\alpha} \sqrt{g}) + \Gamma_{\alpha\beta}^{\beta} \delta(\Gamma_{\mu\nu}^{\alpha} g^{\mu\nu} \sqrt{g}) \\ - \Gamma_{\alpha\beta}^{\beta} \Gamma_{\mu\nu}^{\alpha} \delta(g^{\mu\nu} \sqrt{g})$$

$$= \Gamma_{\mu\nu}^{\alpha} \delta(g^{\mu\nu} \partial_{\alpha} \sqrt{g}) - \Gamma_{\alpha\beta}^{\beta} \delta(\partial_{\nu} (g^{\alpha\nu} \sqrt{g})) \\ - \Gamma_{\alpha\beta}^{\beta} \Gamma_{\mu\nu}^{\alpha} \delta(g^{\mu\nu} \sqrt{g})$$

$$\delta(\Gamma_{\mu\alpha}^{\beta} \Gamma_{\nu\beta}^{\alpha} g^{\mu\nu} \sqrt{g})$$

$$= \partial_{\alpha} (\delta \Gamma_{\mu\alpha}^{\beta}) \Gamma_{\nu\beta}^{\alpha} g^{\mu\nu} \sqrt{g} + \Gamma_{\mu\alpha}^{\beta} \Gamma_{\nu\beta}^{\alpha} \delta(g^{\mu\nu} \sqrt{g})$$

$$= \partial_{\alpha} \delta(\Gamma_{\mu\alpha}^{\beta} g^{\mu\nu} \sqrt{g}) \Gamma_{\nu\beta}^{\alpha} - \Gamma_{\mu\alpha}^{\beta} \Gamma_{\nu\beta}^{\alpha} \delta(g^{\mu\nu} \sqrt{g})$$

$$= -\delta(g^{\nu\beta} \partial_{\alpha} \sqrt{g}) \Gamma_{\nu\beta}^{\alpha} - \Gamma_{\mu\alpha}^{\beta} \Gamma_{\nu\beta}^{\alpha} \delta(g^{\mu\nu} \sqrt{g})$$

$$\Rightarrow \delta \mathcal{L} = \Gamma_{\mu\nu}^{\alpha} \delta \partial_{\alpha} (g^{\mu\nu} \sqrt{g}) - \Gamma_{\alpha\beta}^{\beta} \delta \partial_{\nu} (g^{\alpha\nu} \sqrt{g})$$

$$+ (\Gamma_{\mu\alpha}^{\beta} \Gamma_{\nu\beta}^{\alpha} - \Gamma_{\alpha\beta}^{\beta} \Gamma_{\mu\nu}^{\alpha}) \delta(g^{\mu\nu} \sqrt{g})$$

$$= \partial_{\alpha} (\Gamma_{\mu\nu}^{\alpha} \delta(g^{\mu\nu} \sqrt{g})) - \partial_{\alpha} \Gamma_{\mu\nu}^{\alpha} \delta(g^{\mu\nu} \sqrt{g})$$

$$- \partial_{\nu} (\Gamma_{\alpha\beta}^{\beta} \delta(g^{\alpha\nu} \sqrt{g})) + \partial_{\nu} \Gamma_{\alpha\beta}^{\beta} \delta(g^{\alpha\nu} \sqrt{g})$$

$$+ (\Gamma_{\mu\alpha}^{\beta} \Gamma_{\nu\beta}^{\alpha} - \Gamma_{\alpha\beta}^{\beta} \Gamma_{\mu\nu}^{\alpha}) \delta(g^{\mu\nu} \sqrt{g})$$

$$\Rightarrow \delta I = \delta \int \mathcal{L} d^4x = \int R_{\mu\nu} \delta(g^{\mu\nu} \sqrt{g}) d^4x$$

with the  $\delta g_{\mu\nu}$  arbitrary, the quantities  $\delta(g^{\mu\nu} \sqrt{g})$  are also independent and arbitrary.

$$\delta g^{\mu\nu} = -g^{\mu\alpha} g^{\nu\beta} \delta g_{\alpha\beta}$$

$$\begin{aligned} & \partial_\sigma (g^{\alpha\mu}) g_{\mu\nu} + g^{\alpha\mu} (\partial_\sigma g_{\mu\nu}) \\ &= \partial_\sigma (g^{\alpha\mu} g_{\mu\nu}) \\ &= \partial_\sigma (g^\alpha_\nu) \\ &= 0 \\ \Rightarrow & g_{\mu\nu} (\partial_\sigma g^{\alpha\mu}) = -g^{\alpha\mu} (\partial_\sigma g_{\mu\nu}) \\ \Rightarrow & g^{\beta\gamma} g_{\mu\nu} (\partial_\sigma g^{\alpha\mu}) = -g^{\beta\gamma} g^{\alpha\mu} (\partial_\sigma g_{\mu\nu}) \\ \Rightarrow & \partial_\sigma g^{\alpha\beta} = -g^{\beta\nu} g^{\alpha\mu} (\partial_\sigma g_{\mu\nu}) \\ \Rightarrow & \delta g^{\mu\nu} = -g^{\mu\alpha} g^{\nu\beta} \delta g_{\alpha\beta} \end{aligned}$$

$$\delta \sqrt{-g} = \frac{1}{2} g^{\alpha\beta} \delta g_{\alpha\beta}$$

$$\Rightarrow \delta (g^{\mu\nu} \sqrt{-g}) = - (g^{\mu\alpha} g^{\nu\beta} - \frac{1}{2} g^{\mu\nu} g^{\alpha\beta}) \sqrt{-g} \delta g_{\alpha\beta}$$

$$\Rightarrow \delta I = - \int R_{\mu\nu} (g^{\mu\alpha} g^{\nu\beta} - \frac{1}{2} g^{\mu\nu} g^{\alpha\beta}) \sqrt{-g} \delta g_{\alpha\beta} d^4x$$

$$= - \int (R^{\alpha\beta} - \frac{1}{2} g^{\alpha\beta} R) \sqrt{-g} \delta g_{\alpha\beta} d^4x$$

$$= 0$$

$$\Rightarrow R^{\alpha\beta} - \frac{1}{2} g^{\alpha\beta} R = 0$$

(Einstein's Vacuum Eq.)

The usual expression for the action density of the electromagnetic field is  $\frac{1}{8\pi} (E^2 - H^2)$

If we write it in the 4-D notation of special relativity given in the form

$$\frac{1}{16\pi} F_{\mu\nu} F^{\mu\nu}$$

This leads to the expression

$$I_{em} = \frac{1}{16\pi} \int F_{\mu\nu} F^{\mu\nu} \sqrt{-g} d^4x$$

For the invariant action in general relativity.

$$\tilde{g}_{\mu\nu} = \Omega^2 g_{\mu\nu}$$

$$R = \Omega^2 (\tilde{R} + 6 \tilde{\square} \omega - 6 \tilde{g}^{\mu\nu} \partial_\mu \omega \partial_\nu \omega)$$

$$\omega \equiv \ln \Omega$$

$$\partial_\mu \omega \equiv \frac{\partial \omega}{\partial \tilde{x}^\mu}$$

$$\tilde{\square} \omega \equiv \frac{1}{\sqrt{\tilde{g}}} \partial_\mu (\sqrt{\tilde{g}} \tilde{g}^{\mu\nu} \partial_\nu \omega)$$

$$\Omega^{-2} = F = \exp\left[\sqrt{\frac{2}{3}} k \phi\right]$$

$$\Omega^{-4} = F^2 = \frac{1}{e^{2\sqrt{\frac{2}{3}} k \phi}} = e^{-2\sqrt{\frac{2}{3}} k \phi}$$

$$\Omega^{-8} = e^{-4\sqrt{\frac{2}{3}} k \phi}$$

$$S = \frac{1}{2K^2} \int d^4x \sqrt{-g} f(R) + \int d^4x \sqrt{-g} \left( \frac{-1}{16\pi} F_{\mu\nu} F^{\mu\nu} \right) + \int d^4x \sqrt{-g} \mathcal{L}_M(g_{\mu\nu}, \Psi_M)$$

$\uparrow$   
 $K = \sqrt{8\pi G}$

$$= \int d^4x \sqrt{-g} \left( \frac{1}{2K^2} FR - \frac{FR-f}{2K^2} \right) + \int d^4x \sqrt{-g} \left( \frac{-1}{16\pi} F_{\mu\nu} F^{\mu\nu} \right) + \int d^4x \sqrt{-g} \mathcal{L}_M(g_{\mu\nu}, \Psi_M)$$

$\equiv$   
 $U$

$$= \int d^4x \sqrt{-g} \left( \frac{1}{2K^2} FR - U \right) + \int d^4x \sqrt{-g} \left( \frac{-1}{16} g_{\alpha\mu} g_{\beta\nu} F^{\alpha\beta} F^{\mu\nu} \right) + \int d^4x \sqrt{-g} \mathcal{L}_M(g_{\mu\nu}, \Psi_M)$$

$$= \int d^4x \sqrt{-g} \Omega^{-4} \left( \frac{1}{2K^2} F(\Omega^2 \tilde{R} + 6 \tilde{\nabla}^\mu \omega - 6 \tilde{g}^{\mu\nu} \partial_\mu \omega \partial_\nu \omega) - U \right)$$

$$+ \int d^4x \sqrt{-g} \Omega^{-8} \left( \frac{-1}{16} \tilde{g}_{\alpha\mu} \tilde{g}_{\beta\nu} F^{\alpha\beta} F^{\mu\nu} \right)$$

$$+ \int d^4x \sqrt{-g} \Omega^{-4} \mathcal{L}_M(\Omega^{-2} \tilde{g}_{\mu\nu}, \Psi_M)$$

$$\left[ \begin{array}{l} \Omega^2 = F \\ K\phi = \sqrt{\frac{3}{2}} \ln F \quad \ln \Omega^2 = \ln F \\ \ln \Omega = \frac{1}{2} \ln F \\ \ln F = \frac{\sqrt{2}}{\sqrt{3}} K\phi \end{array} \right. \quad \left. \begin{array}{l} \ln \Omega = \frac{1}{2} \ln F \\ = \frac{1}{2} \left( \frac{\sqrt{2}}{\sqrt{3}} K\phi \right) = \frac{1}{\sqrt{6}} K\phi = \omega \end{array} \right.$$

$$= \int d^4x \sqrt{-g} \left[ \frac{1}{2K^2} \tilde{R} - \frac{1}{2} \tilde{g}^{\mu\nu} \partial_\mu \phi \partial_\nu \phi - V(\phi) \right]$$

$$- \frac{1}{16} \int d^4x \sqrt{-g} \left[ F^{-4}(\phi) \tilde{g}_{\alpha\mu} \tilde{g}_{\beta\nu} F^{\alpha\beta} F^{\mu\nu} \right]$$

$$+ \int d^4x \sqrt{-g} \left[ F^{-2}(\phi) \mathcal{L}_M(F^{-1}(\phi) \tilde{g}_{\mu\nu}, \Psi_M) \right]$$